



## Three Dimensional Casson nanofluid Flow with Convective Boundary Layer via Stretching Sheet

G. Radha<sup>1</sup>, B. Reddappa<sup>2</sup>, Nainaru Tarakaramu<sup>3\*,4\*</sup>, V.K. Somasekhar Srinivas<sup>5,6</sup>, Sivajothi Ramalingam<sup>7</sup>, N. Mallikarjuna Reddy<sup>8</sup>, Kotte Amaranadha Reddy<sup>9</sup>, K. Bhagya Lakshmi<sup>10</sup>

<sup>1</sup> Associate Professor of Mathematics, Department of Humanities and Science, KSRRM College of Engineering (Autonomous), Kadapa-516003, Andrapradesh, India

<sup>2,9</sup> Department of Mathematics, Kalasalingam Academy of Research and Education (Deemed to be University), Krishnankoil, Srivilliputhur, Tamilnadu-626126, India

<sup>3\*,5</sup> Department of Mathematics, School of Liberal Arts and Sciences, Mohan Babu University, Sree Sainath Nagar, Tirupati-517102, Andrapradesh, India

<sup>4\*,6</sup> Department of Mathematics, Basic Science and Humanities, Sree Vidyanikethan Engineering College, Sree Sainath Nagar, Tirupati-517102, Andrapradesh, India

<sup>7</sup> Department of Management, R L Institute of Management Studies (A Unit of Subbalakshmi Lakshmi Pathy College of Science) Madurai-625022, Tamil Nadu, India

<sup>8</sup> RJS FIRST GRADE COLLEGE, Koramangala, Bangaluru-560034, Karnataka, India

<sup>10</sup> Department of Mathematics, Humanities and Basic Sciences, Annamacharya Institute of Technology & Sciences, Tirupati, Andhra Pradesh, India

\*Corresponding author's Email: [nainaru143@gmail.com](mailto:nainaru143@gmail.com)

Article History	Abstract
Received: 06 June 2023 Revised: 05 Sept 2023 Accepted: 20 Oct 2023	<p>The present work examined Casson nanofluid in a three-dimensional boundary layer motion via stretching sheet. The study focuses on analyzing the behavior of a Casson nanofluid, which is one type of non-Newtonian fluid. The study appears to involve solving partial differential equations related to fluid flow, heat transfer, and mass transfer. These PDEs are transformed into ordinary differential equations using standard similarity variables. To solve the ODEs, the researchers employ the Runge-Kutta-Fehlberg (R-K-F) 4<sup>th</sup> order iterative scheme. It appears that higher values of the Biot number can significantly affect the temperature and concentration profiles in the Casson liquid flow.</p> <p><b>Keywords:</b> Casson Fluid, Nanofluid, Stretching Sheet, Convective.</p>
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### 1. Introduction

The study of nanofluids has drawn a lot of interest in the field of fluid mechanics and heat transfer, primarily owing to their unique thermal characteristics and potential usage in various engineering and industrial processes. Firas et al. [1] have investigated the steady laminar 2-D incompressible MHD natural convection flow that occurs around a solid sphere submerged in a Casson nanofluid. Shah et al. [2] investigated entropy optimization in electrically conducting Casson nanofluid flowing over a stretchable nonlinear surface. Jamshed et al. [3] examined role of solar thermal transport in Casson nanofluid with slip condition. Adebowale Martins et al. [4] evaluated the second-order velocity slip and heat transfer caused by nanofluid along with non-Darcian Casson flow over a permeable stretching surface. Abo-Dahab et al. [5] studied the viscoelastic fluid flow over a nonlinearly stretched surface. Sahoo and Nandkeolyar [6] examined entropy production in a three-dimensional Casson nanofluid flow that conducts electricity. Akaje and Olajuwon [7] examined the impact of nonlinear radiative heat on species heat transfer while taking Thompson and Troian boundary conditions. Satya Narayana et al. [8]

investigated the fluid flow over an internally heated and horizontally stretched surface. Akinshilo et al. [9] examined the non-Newtonian Casson nanofluid flow around a small needle. Mahanta et al. [10] focused on the real-world uses of nanoparticles in non-Newtonian base fluids, particularly in the fields of energy conversion and heat generation.

Suresh Kumar et al. [11] studied the behavior of Casson nanofluid flow via an exponentially stretching surface, with activation energy, thermal radiation, Brownian motion, and thermophoresis. Taj and Salahuddin [12] examined the effect of viscous dissipation and nonlinear radiation in three-dimensional Williamson fluid flow via exponentially stretching surface. Ibrahim and Fawzy [13] investigated the effects of Arrhenius activation energy in a rotating flow of Casson nanofluid. Wang et al. [14] studied the impact of activation energy on 3D nanofluid motion over a stretching surface with slip conditions. Showkat Ahmad et al. [15] encountered the non-Newtonian Casson fluid's magnetohydrodynamic flow over a stratified extending sheet Jana Reddy et al. [16] investigated the Casson nanofluid flow characteristics over a permeable, non-linearly stretched sheet. Wang et al. [17] examined the effects of activation energy and chemical reactions in non-Newtonian liquid flow over a stretching sheet. Suresh Kumar et al. [18] studied the Brownian and thermophoresis effects with activation energy and nonlinear thermal radiation on mixed convection Casson fluid flow over a vertical cone. Tarakaramu et al. [19] have studied the non-Newtonian nanofluid flow in three dimensions over a porous stretching sheet. Zeeshan et al. [20] have evaluated various flow scenarios in applied sciences, particularly in areas like nuclear reactor cooling. Khan et al. [21] have examined the generation of entropy during the radiative spinning motion of a Casson nanofluid. Alrehili et al. [22] contributed many engineering applications by investigating the behavior of dissipative Carreau nanofluids flowing. Abideen et al. [23] have investigated heat transfer in a radiative nanofluid flowing over a curved surface. Recently, some of authors [24-26] developed three dimensional Casson nanofluid motion via stretching sheet.

## 2. Materials And Methods

### Mathematical Analysis

The mathematical model as consider by using following aspects such as:

- Consider the 3D convective incompressible couple stress non-Newtonian Casson liquid motion via with thermal radiation.
- We considered the liquid motion through  $x^*$ ,  $y^*$  directions.
- The fluid motion considers at stretching sheet in  $z^*$  is vanishes.
- The velocity components of axial and transverse direction  $u = U_w^*(x^*) = a_1 x^*$ ,  $v = V_w^*(x^*) = b_1 y^*$  as shown in **Fig.1**.

The rheological equation of Casson liquid motion has to be consider as follows

$$\tau_{ij} = \begin{cases} \left( 2\mu_0^* + \frac{2p_y^*}{\sqrt{2\pi^*}} \right) e_{ij}, & \text{if } \pi^* \geq \pi_1^* \\ \left( 2\mu_0^* + \frac{2p_y^*}{\sqrt{2\pi_1^*}} \right) e_{ij}, & \text{if } \pi^* < \pi_1^* \end{cases} \tag{1}$$

Where,  $p_y^* = e_{ij}e_{ij}$  and  $\beta = \mu_B^* \sqrt{2\pi_1^*} / p_y^*$ .

The fundamental equations of continuity, heat and concentration Eqs can be formulated by using eq. (1)

$$\frac{\partial u}{\partial x^*} + \frac{\partial v}{\partial y^*} + \frac{\partial w}{\partial z^*} = 0 \quad (2)$$

$$u \frac{\partial u}{\partial x^*} + v \frac{\partial u}{\partial y^*} + w \frac{\partial u}{\partial z^*} = \nu^* (1 + \beta^{-1}) \frac{\partial^2 u}{\partial (z^*)^2} \quad (3)$$

$$u \frac{\partial v}{\partial x^*} + v \frac{\partial v}{\partial y^*} + w \frac{\partial v}{\partial z^*} = \nu^* (1 + \beta^{-1}) \frac{\partial^2 v}{\partial (z^*)^2} \quad (4) (4)$$

$$u \frac{\partial T^*}{\partial x^*} + v \frac{\partial T^*}{\partial y^*} + w \frac{\partial T^*}{\partial z^*} = \alpha \frac{\partial^2 T^*}{\partial z^{*2}} + \tau \left( D_B \frac{\partial T^*}{\partial z^*} \frac{\partial C}{\partial z^*} + \frac{D_T}{T_\infty} \left( \frac{\partial T^*}{\partial z^*} \right)^2 \right) \quad (5)$$

$$u \frac{\partial C}{\partial x^*} + v \frac{\partial C}{\partial y^*} + w \frac{\partial C}{\partial z^*} = \left( D_B \frac{\partial^2 C}{\partial (z^*)^2} + \frac{D_T}{T_\infty} \left( \frac{\partial T^*}{\partial (z^*)^2} \right) \right) \quad (6)$$

The present relevant model boundary conditions as

$$\begin{aligned} z^* = 0 \quad \text{at} \quad u_1 = a_1 x^* \quad v = b_1 y^* \quad w_1 = 0, \quad -k \frac{\partial T^*}{\partial z^*} = H_1 (T_f^* - T^*) \quad -D^* \left( \frac{\partial C}{\partial z^*} \right) = H_2 (C_f - C) \\ z^* \rightarrow \infty \quad \text{as} \quad u_1 \rightarrow 0 \quad v_1 \rightarrow 0, \quad u_1' \rightarrow 0, \quad v_1' \rightarrow 0, \quad T^* \rightarrow T_\infty^*, \quad C \rightarrow C_\infty \end{aligned} \quad (7)$$

The similarity transformations as below

$$\left. \begin{aligned} \eta_1 = \sqrt{\frac{a_1}{\nu_f^*}} z^*, \quad u = a_1 x^* f'(\eta_1), \quad v = a_1 y^* g'(\eta_1), \quad w = -\sqrt{a_1 \nu^*} (f(\eta_1) + g(\eta_1)) \\ \theta(\eta_1) = \frac{T - T_\infty}{T_w - T_\infty}, \quad \phi(\eta_1) = \frac{C - C_\infty}{C_w - C_\infty} \end{aligned} \right\} \quad (8)$$

Using above **Eq. (8)**, we are converting **Eq. (3)-(6)** into below format

$$f''' = \left[ \frac{1}{(1 + \beta/\beta)} \right] [-f''(f + g) + 2f'(f' + g')] \quad (9)$$

$$g''' = \left[ \frac{1}{(1 + \beta/\beta)} \right] [-g''(f + g) + 2g'(f' + g')] \quad (10)$$

$$\theta'' = -\text{Pr} \left[ -(f + g)\theta' - N_b \theta' \phi' - N_t (\theta')^2 \right] \quad (11)$$

$$\phi'' = -\text{Le} \left( (f + g)\phi' - \left( \frac{N_t}{N_b} \right) \text{Pr} \left[ -(f + g)\theta' - N_b \theta' \phi' - N_t (\theta')^2 \right] \right) \quad (12)$$

Corresponding B.Cs. as below

$$\left. \begin{aligned} f = 0, \quad g = 0, \quad f' = 1, \quad g' = 1, \quad \theta' = -Bi_1(1 - \theta), \quad \phi' = -Bi_2(1 - \phi) \quad \text{at} \quad \eta_1 = 0 \\ f' \rightarrow 0, \quad g' \rightarrow 0, \quad \theta \rightarrow 0, \quad \phi \rightarrow 0, \quad \text{as} \quad \eta_1 \rightarrow \infty \end{aligned} \right\} \quad (13)$$

Moreover, the skin-friction coefficient and Nusselt number are below

$$\text{Re}_x^{0.5} C_{fx} = \left(1 + \frac{1}{\beta}\right) f''(0), \quad \text{Re}_x^{0.5} C_{fy} = \left(1 + \frac{1}{\beta}\right) g''(0) \quad (14).$$

### 3. Results and Discussion

#### Analysis of Phytochemicals and Percentage yield

**Fig. 2** predicts the  $\beta$  (Casson Fluid Parameter) on  $f'(\eta)$ . These observations suggest that the Casson fluid parameter plays a critical role in determining both the velocity profile within the boundary layer. Increasing  $\beta$  results in reduced velocity.

**Fig. 3** presented  $Pr$  (Prandtl number) on  $\theta(\eta)$ . This behaviour is consistent with the physical interpretation of the Prandtl number. A higher Prandtl number implies that heat is conducted less efficiently within the fluid compared to the momentum transfer (velocity). This can have significant implications for heat transfer processes, such as conduction, convection, and boundary layer phenomena.

**Fig. 4(a)-4(b)** presented the  $N_b$  (Brownian Motion Parameter) on  $\theta(\eta)$ ,  $\phi(\eta)$ . We have described characteristic of the influence of Brownian motion on nanoparticle dispersion and transport within a fluid. When  $N_b$  is higher, it implies that Brownian motion becomes more significant, causing nanoparticles to move more erratically within the fluid. This increased movement can lead to enhanced mixing and heat transfer, which, in turn, results in a thicker thermal boundary layer and higher temperatures within the nanofluid.

In summary of  $N_t$  (Thermophoresis Parameter) exhibited on  $\theta(\eta)$ ,  $\phi(\eta)$  as shown in **Figs. 4(a)-4(b)**, respectively. The rise in fluid temperature and concentration with an increase in the thermophoresis parameter  $N_t$  is linked to the behavior of nanoparticles within the fluid.

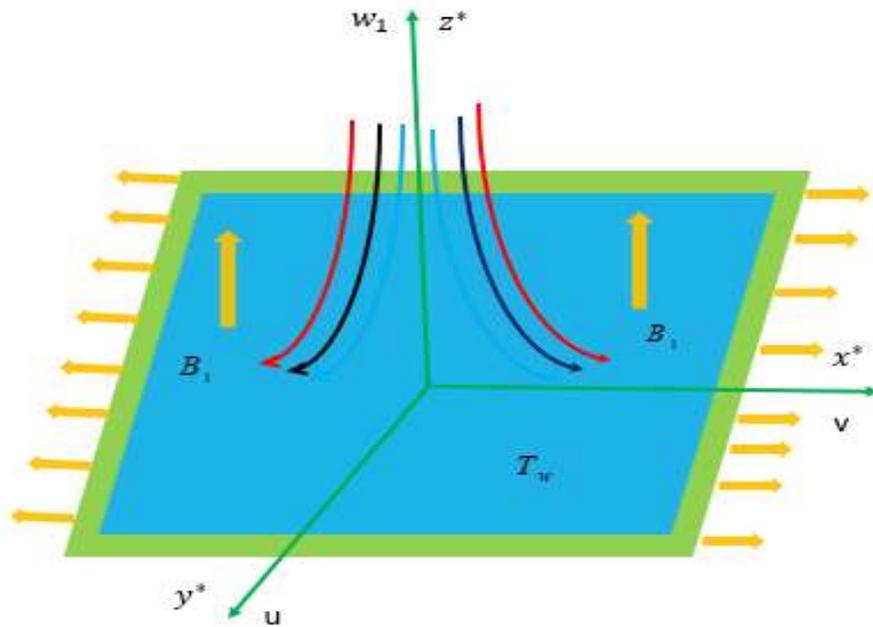
**Fig. 5** shows characteristics of Lewis number on  $\phi(\eta)$ . These observations indicate that the Lewis number plays a significant role in determining the behavior of nanoparticle concentration profile in the presence of Brownian diffusion. A higher Lewis number implies that thermal diffusion is more dominant relative to molecular diffusion, leading to decreased Brownian diffusion and, consequently, a thinner concentration boundary layer and reduced nanoparticle concentration near the sheet.

**Table. 1** presented the variation of Casson Fluid Parameter on Skin friction coefficients

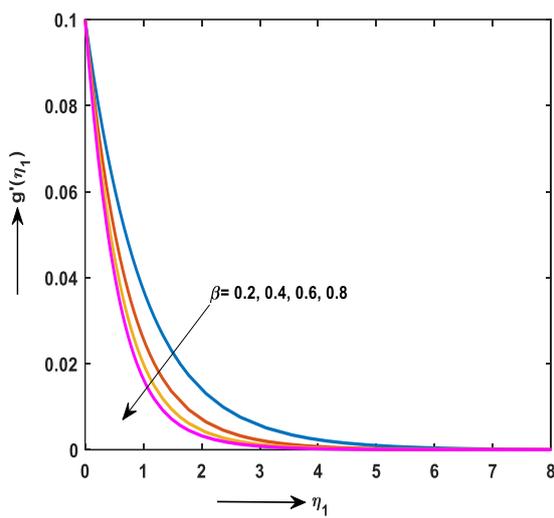
#### Conclusions:

The main contribution results as presented as following:

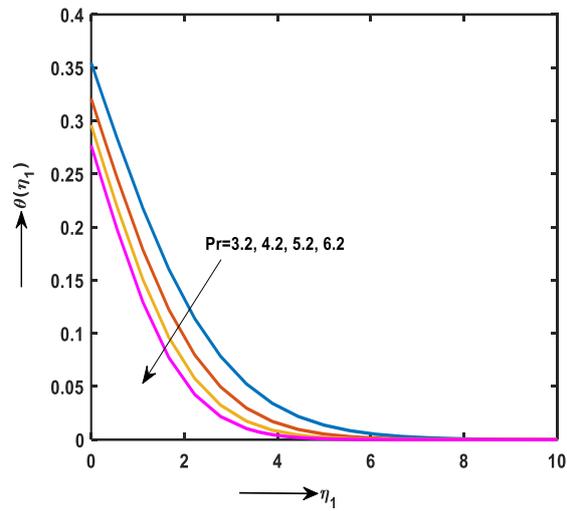
- ❖ The velocity of Casson fluid is declined with large statistical values of Casson liquid parameter.
- ❖ The temperature and concentration are high for enlarge values of Biot number of temperature and concentration.



**Fig. 1** Flow Geometry



**Fig. 2** Influence of  $\beta$  on  $g'(\eta_1)$



**Fig. 3** Influence of  $Pr$  on  $\theta(\eta)$

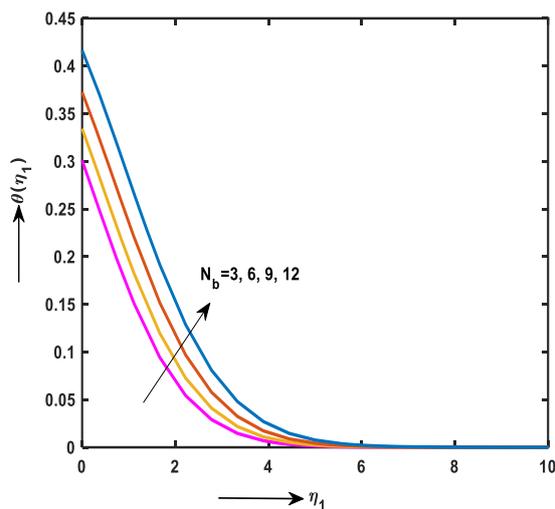


Fig. 3 Influence of  $N_b$  on  $\theta(\eta)$

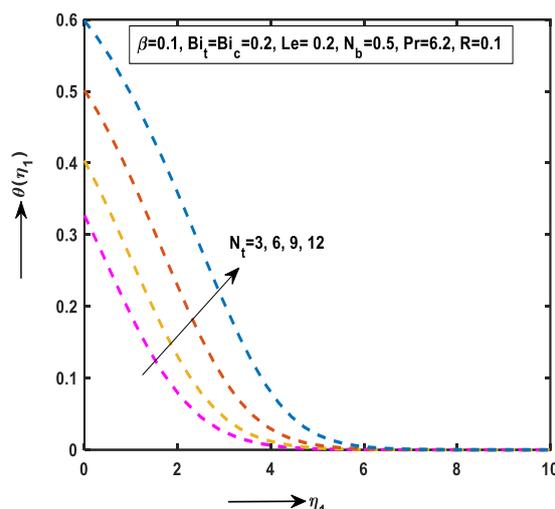


Fig. 4(a) Influence of  $N_b$  on  $\theta(\eta)$

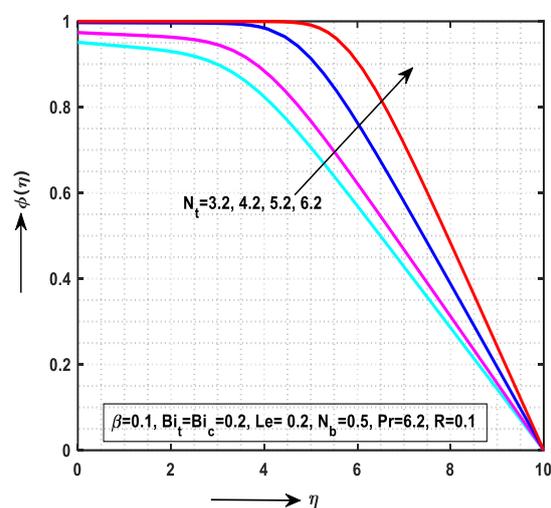


Fig. 4(b) Influence of  $N_b$  on  $\phi(\eta)$

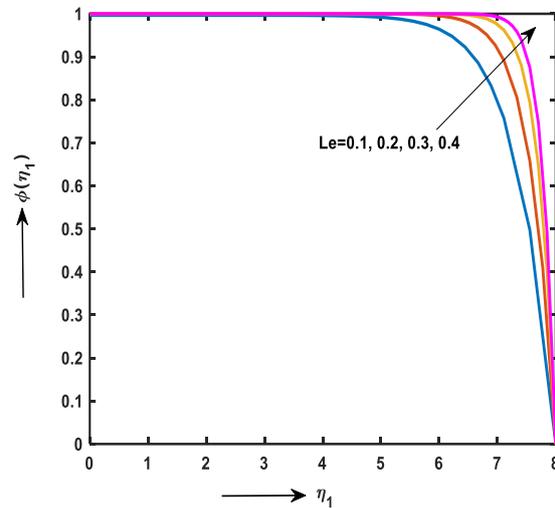


Fig. 5 Influence of  $Le$  on  $\phi(\eta)$

Table 1. Variation of the Casson Fluid Parameter on Skin friction coefficients:

$\beta$	$-\left(\frac{1}{1+\beta}\right)f''(\eta)$	$-\left(\frac{1}{1+\beta}\right)g''(\eta)$
$\infty$	1.2105	0.00945
0.5	1.2212	0.12891
1	2.1221	0.32548
1.5	2.2234	0.29358
2	3.0023	0.32546
2.5	3.1589	0.33564
3	3.2558	0.45689
3.5	3.2258	0.44512
4	3.3348	0.45891
4.5	3.4568	0.48901
5	3.5891	0.51254

Nomenclature	
$a_1, b_1$ Constants	$T_\infty$ Ambient fluid temperature
$Nu_x$ Nusselt number	$U_\infty$ Free stream velocity
$u, v, w$ Velocity components along $x^*, y^*, z^*$	$U_w, V_w$ Stretching velocities
$C$ Nanoparticle volume fraction	$W_c$ Maximum cell swimming speed
$C_p$ Specific heat constant ( $kJ / kg K$ )	$\lambda_i$ Slip factors $= N_i \sqrt{a_1 \mu^*}$
$C_f$ Skin friction coefficient	Greek symbols
$C_\infty$ Uniform ambient concentration ( $Kg m^{-3}$ )	$\rho_f$ Fluid density
$C_w$ Nanoparticle concentration ( $Kg m^{-3}$ )	$(\rho c)_p$ Heat capacity of the nanoparticle
$D_n$ Diffusivity of microorganisms	$(\rho C)_{bf}$ Base fluid
$D_B$ Brownian diffusion	$\rho_f$ Fluid density
$D_T$ Thermophoresis diffusion ( $m^2 .s^{-1}$ )	$(\rho c)_f$ Heat capacity of the field ( $kJ kg^{-1}$ )
$f'$ Dimensionless velocity	$\phi$ Dimensionless concentration
$f$ Dimensionless stream function	$\eta_1$ Similarity variable
$h_f$ Heat transfer coefficient	$\mu^*$ Dynamic viscosity ( $Pa.s^{-1}$ )
$k^*$ Mean absorption coefficient	$\theta$ Dimensionless temperature
$Le$ Lewis number $\frac{\alpha_m}{D_B}$	$\rho_f$ Fluid density ( $Kg .m^{-3}$ )
$N_t$ Thermophoresis parameter $= \frac{\tau^* D_T (T_w - T_\infty)}{\alpha_m T_\infty}$	$Bi_t$ Surface Convection Parameter $\frac{h_f}{k} (\sqrt{\nu/a})$
$N_b$ Brownian motion coefficient $= \frac{\tau^* D_B (C_w - C_\infty)}{\alpha_m}$	$Bi_c$ Surface Convection Parameter $\frac{h_s}{D_B} (\sqrt{\nu/a})$
$N_1, N_2$ Slip coefficients in x and y	$\tau^*$ Ratio of the nanoparticle to fluid $\frac{(\rho c)_p}{(\rho c)_f}$
$Pr$ Prandtl number $= \frac{\nu^*}{\alpha_m}$	$\nu^*$ Kinematic viscosity $= \frac{\mu^*}{\rho_f} (m^2 s^{-1})$
$R$ Thermal Radiation $= \frac{16\sigma^* T_\infty^3}{3k^* K^*}$	Subscripts
$T_f$ Temperature of hot fluid	$\infty$ condition at free stream
$T$ Fluid temperature	$w$ wall mass transfer velocity ( $m s^{-1}$ )

$Re_x$ Reynolds number	
$Sc$ Schmidt number	$= \frac{v^*}{D_n}$
<i>Abbreviations</i>	
HT Heat Transfer	TC Thermal Conductivity
TR Thermal Radiation	NPs Nanoparticles
NFs Nanofluids	SS Stretching Sheet
CF Casson Fluid	

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